# 250. The Interaction of Non-bonding Orbitals in Dicarbonyls: *Ab Initio* Results on Glyoxal

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Zusammenfassung. Die energetische Lage der kanonischen Molekelorbitale von Glyoxal in der trans- und cis-Konformation wurde durch ab initio-Rechnungen bestimmt. Die Wechselwirkung der «nichtbindenden» Orbitale dieser Molekel wird diskutiert.

Until recently spectroscopists have used simple 2p-functions localized on the oxygen atoms to describe the non-bonding orbitals (*n*-orbitals) in dicarbonyl compounds [1]. Molecular orbital calculations [2] of the UV. and circular dichroism spectra by semi-empirical methods including all valency electrons indicate that such a description is, even qualitatively, insufficient. Experimental evidence from photo-electron spectroscopy [3], as well as from electrochroism measurements [4], supports this. However, semi-empirical calculations are in general subject to the arbitrariness

	· · · · · · · · · · · · · · · · · · ·	trans (C <sub>2h</sub> ) b)	$cis(\mathbf{C}_{2v})\mathbf{b})$
MO		$-20.6415 (1 a_g)$	$-20.6464(1a_1)$
		$-20.6415 (1 b_g)$	$-20.6464(1 b_1)$
		$-11.5314(2a_g)$	$-11.5250(2a_1)$
		$-11.5308(2b_g)$	$-11.5244(2b_1)$
		$-1.4876(3 a_g)$	$-1.4847(3a_1)$
		$-1.4620(3b_g)$	$-1.4542(3b_1)$
		$-1.0056 (4 a_g)$	$-1.0133(4a_1)$
		$-0.8689 (4 b_g)$	$- 0.8385 (4b_1)$
		$-0.7604(5a_g)$	$- 0.7946(5a_1)$
		$-0.7113(5b_g)$	$- 0.6826(6a_1)$
		$-0.6632(6a_g)$	$- 0.6721(5b_1)$
	π:	$- 0.6317 (1 a_u)$	$- 0.6269 (1 b_2)$
	n:	$-0.5635(6b_g)$	$-0.5554(6b_1)$
	$\pi$ :	$- 0.5458 (1 b_u)$	$- 0.5404 (1 a_2)$
	<i>n</i> :	$-$ 0.4768 (7 $a_g$ )	$- 0.4775 (7 a_1)$
$\overline{E_T}$		– 226.24765	- 226.24288

Table 1. Orbital energies in trans- and cis-Glyoxala)

a) Atomic units.

b) For the cis-form as yet unknown, we have chosen the standardized geometry: C-H: 1.08 Å, C-O: 1.22 Å, C-C: 1.5 Å, and <OCH: 120°.</p>

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of the choice of empirical parameters. We have therefore effected an *ab initio* calculation<sup>2</sup>) on glyoxal, the simplest  $\alpha$ -dicarbonyl compound. Such a calculation should provide an unambiguous decision concerning the MO level ordering as well as the energy gap between the non-bonding orbitals.

Ab initio molecular wave-functions have been constructed with Gaussian lobefunctions as a basis set, with an accuracy equivalent to double zeta Slater-type orbitals

		$cis[C_{2v}]$		trans [C <sub>2</sub> h]	
		symm. (7 <i>a</i> <sub>1</sub> )	anti- symm. (6 <i>b</i> 1)	symm. (7 <i>a</i> <sub>g</sub> )	anti- symm. $(6b_g)$
$\varphi_{1s}$	(C1)	0.0157	- 0.0334	0.0158	- 0.0345
\$\$15-25	(C1)	0.0007	-0.0017	0.0007	-0.0018
$\varphi_{2s}$	(C1)	-0.0433	0.2167	0.0458	0.2217
$\varphi_{2px}$	(C1)	0.3351	- 0.0980	0.3324	- 0.0908
$\varphi_{2py}$	(C1)	-0.1191	-0.1585	-0.1242	- 0.1551
$\varphi_{2p_z}$	(C1)	0.0000	0.0000	0.0000	0.0000
$\varphi_{1s}$	(C2)	0.0157	0.0334	0.0158	0.0345
$\varphi_{1s-2s}$	(C2)	0.0007	0.0017	0.0007	0.0018
$\varphi_{2s}$	(C2)	- 0.0433	-0.2167	-0.0458	-0.2217
$\varphi_{2p_x}$	(C2)	- 0.3351	- 0.0980	-0.3324	- 0.0908
$\varphi_{2py}$	(C2)	- 0.1191	0.1585	0.1242	- 0.1551
$\varphi_{2p_z}$	(C2)	0.0000	0.0000	0.0000	0.0000
$\varphi_{1s}$	(O1)	0.0014	- 0.0091	0.0013	- 0.0080
\$\$\varphi_{1s-2s}\$	(O1)	0.0001	- 0.0005	0.0001	- 0.0004
$\varphi_{2s}$	(O1)	- 0.0080	0.0627	- 0.0084	0.0593
$\varphi_{2p_X}$	(O1)	- 0.5790	0.4941	-0.5871	- 0.4896
$\varphi_{2py}$	(O1)	0.2734	0.4505	0.2700	0.4508
$\varphi_{2p_z}$	(O1)	0.0000	0.0000	0.0000	0.0000
$\varphi_{1s}$	(O2)	0.0014	0.0091	0.0013	0.0080
\	(O2)	0.0001	0.0005	0.0001	0.0004
$\varphi_{25}$	(O2)	-0.0080	0.0627	- 0.0084	- 0.0593
$\varphi_{2p_{X}}$	(O2)	0.5790	0.4941	0.5871	- 0.4896
$\varphi_{2py}$	(O2)	0.2734	0.4505	- 0.2700	0.4508
$\varphi_{2p_z}$	(O2)	0.0000	0.0000	0.0000	0.0000
$\varphi_{H1}$		0.2156	0.1619	0.2121	0.1619
$\varphi_{\mathbf{H2}}$		0.2156	- 0.1619	0.2121	- 0.1619
$\varepsilon_i b$		-0.4775	- 0.5554	0.4768	- 0.5635

Table 2. n-Orbitals in Glyoxal<sup>a</sup>)

<sup>a)</sup>  $\varphi_{1s}, \varphi_{1s-2s}, \varphi_{2s}$  are the Gaussian group orbital basis set which represent the 1s and 2s Hartree-Fock atomic orbitals. For example the carbon 1s atomic orbital is composed almost entirely of  $\varphi_{1s}$  and the carbon 2s almost entirely of  $\varphi_{2s}, \varphi_{1s-2s}$  is intermediate. It is simply an additional basis function that makes possible a more perfect match of the atomic Hartree-Fock solution. The molecule lies in the xy-plane, C—C bond being the x-axis.

b) Atomic units,  $\varepsilon_i$  is the MO energy.

<sup>2</sup>) Details on the wave functions and the *cis-trans*-isomerization process will be published elsewhere; *T.-K. Ha*, unpublished results.

[5]. The wave-functions employ 10 spherical *Gaussians* for the 1s and 2s orbitals and 5 pairs of Gaussians for each  $2p_x$ ,  $2p_y$ , or  $2p_z$  orbital of the carbon and oxygen atoms. Each hydrogen 1s orbital is represented by 5 spherical *Gaussians*.

Tab. 1 shows that the four highest occupied MO's are the two *n*- and the two  $\pi$ -orbitals, the sequence being  $n-\pi-n-\pi$ . We obtain for both *cis*- and *trans*-glyoxal an antisymmetric combination for the *n*-orbital below a symmetric one. The energy gap amounts to 2.36 eV in the *trans*- and 2.12 eV in the *cis*-form. The experimental value from photoelectron spectra is 1.6 eV [3].

The energy gap is nearly independent of the dihedral angle between the two carbonyl groups. A direct interaction between the 2p-orbitals on the oxygen atoms would place the symmetric combination below the antisymmetric one. The distance between the oxygen atoms being rather large in the *trans*-conformation (3.4 Å), we may consider the decrease of the energy gap by 0.24 eV in the *cis*-form to represent a rough measure of the "through space" interaction in *cis*-glyoxal. The value of 0.24 eV seems, however, rather large and considerably exceeds the values estimated from *Slater*-type orbitals [1].

The eigenvectors of the two *n*-orbitals are shown in Tab. 2. The values given provide means for a qualitative understanding of the large energy gap between the two *n*-orbitals: Contour diagrams<sup>2</sup>) of the electronic charge distribution for these orbitals show that, in contrast to the higher lying orbital, the lower one is devoid of nodes bisecting the C-O bonds.

Based on experimental findings we thus show definitely that the mechanism of the interactions of non-bonding orbitals in dicarbonyl compounds is "through bond" rather than "through space".

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